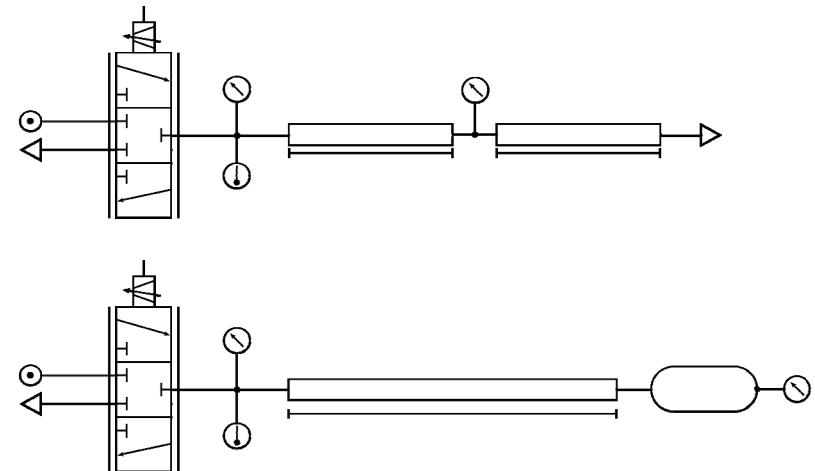




# Modelling of a Long Pneumatic Transmission Line: Models of Successively Decreasing Complexity and their Experimental Validation

Ampère Lab

Lyon, 27th of October 2016



# Overview

## Motivation

### Purpose of a model

- Analysis
- Feedforward control
- Feedback control
- Optimization

### Problems

- Which model is well-suited to describe a pneumatic transmission line?
- „Le simple est toujours faux. Ce qui ne l'est pas est inutilisable.“ Paul Valéry (1937)



$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + (\rho v^2 + p)_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

$$(\rho e)_t + (v(\rho e + p))_z = \frac{1}{A} 2\pi \alpha r_i (T_0 - T) + f_{\text{comp}} \frac{\rho v^2 |v|}{2D}$$

Simplifications

$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + a_{\text{iso}}^2 \rho_z = -k_{\text{fric}} \frac{32\eta_0}{D^2} \frac{1}{\rho_0} \rho v$$

# Overview

## Agenda

### Overview

- Motivation
- Agenda

### One-dimensional compressible flow

- Conservation laws
- Friction and heat transfer
- Successive model simplifications

### Simulation and measurement results

- Scenario 1 – Flow into ambient air
- Scenario 2 – Flow into a terminating volume

### Summary and outlook

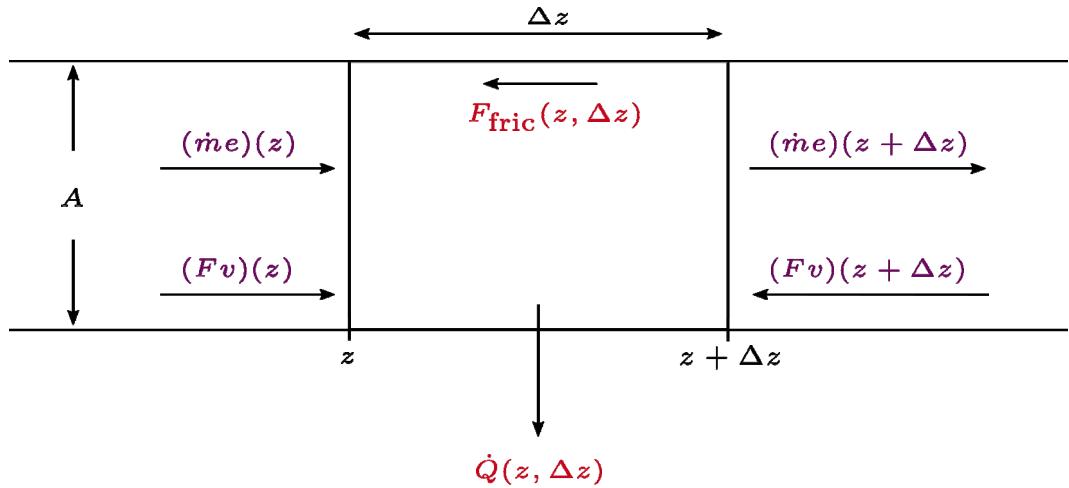
### References

# One-dimensional compressible flow

## Conservation laws

### Euler equations

- Conservation of mass  
 $\rho_t + (\rho v)_z = 0$
- Conservation of momentum  
 $(\rho v)_t + (\rho v^2 + p)_z = 0$
- Conservation of energy  
 $(\rho e)_t + (v(\rho e + p))_z = 0$



### Augmented Euler equations

- Conservation of mass  
 $\rho_t + (\rho v)_z = 0$
- Conservation of momentum  
 $(\rho v)_t + (\rho v^2 + p)_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$
- Conservation of energy  
 $(\rho e)_t + (v(\rho e + p))_z = \frac{1}{A} 2\pi \alpha r_i (T_0 - T) + f_{\text{comp}} \frac{\rho v^2 |v|}{2D}$

# One-dimensional compressible flow

## Friction and heat transfer

### Augmented Euler equations

- Conservation of mass

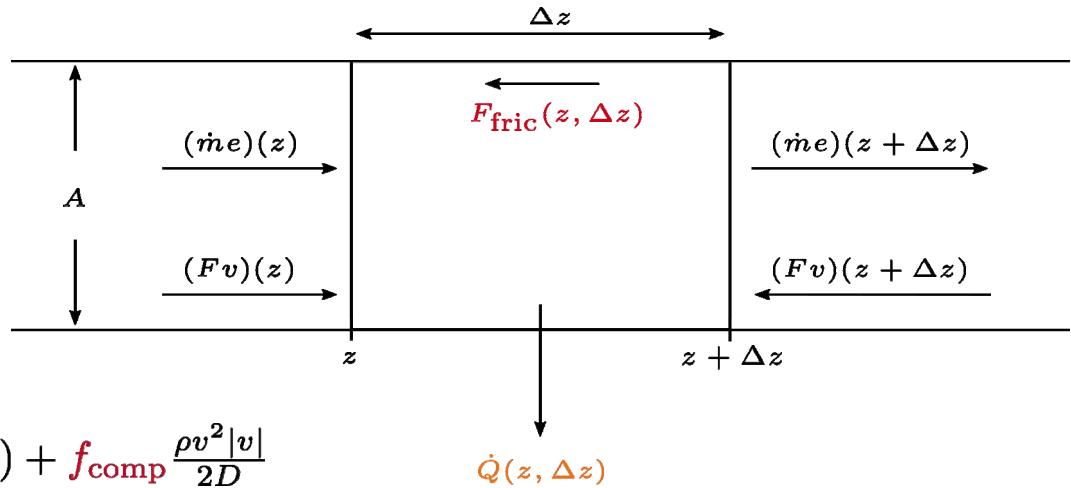
$$\rho_t + (\rho v)_z = 0$$

- Conservation of momentum

$$(\rho v)_t + (\rho v^2 + p)_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

- Conservation of energy

$$(\rho e)_t + (v(\rho e + p))_z = \frac{1}{A} 2\pi \alpha r_i (T_0 - T) + f_{\text{comp}} \frac{\rho v^2 |v|}{2D}$$



### Empirical correlations

- Compressible friction factor

$$f_{\text{comp}} = f \left(1 + \frac{\gamma-1}{2} Ma^2\right)^{-0.47}$$

- Laminar flow

$$f = \frac{64}{Re}$$

- Turbulent flow (Haaland)

$$\frac{1}{\sqrt{f}} = -1.8 \log \left[ \left( \frac{\varepsilon/D}{3.7} \right)^{1.11} + \frac{6.9}{Re} \right]$$

- Heat transfer coefficient (Gnielinski)

$$\alpha = \frac{\lambda_f}{D} \frac{\frac{f_{\text{comp}}}{8} (Re - 1000) Pr}{1 + 12.7 \left( \frac{f_{\text{comp}}}{8} \right)^{\frac{1}{2}} \left( Pr^{\frac{2}{3}} - 1 \right)}$$

# One-dimensional compressible flow

## Successive model simplifications

### Model 1

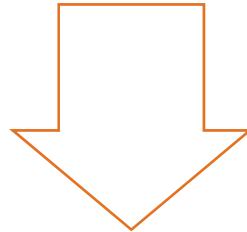
$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + (\rho v^2 + p)_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

$$(\rho e)_t + (v(\rho e + p))_z = \frac{1}{A} 2\pi \alpha r_i (T_0 - T) + f_{\text{comp}} \frac{\rho v^2 |v|}{2D}$$

### Assumptions

- No diffusion
- One-dimensional flow
- No gravity
- Fluid is a polytropic, ideal gas



### Model 2

$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + (\rho v^2 + a_{\text{iso}}^2 \rho)_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}, \quad a_{\text{iso}} = \sqrt{\gamma R_s T_0}$$

### Additional Assumption

- Flow is isothermal
  - Valid if pressure changes are slow compared to the time needed to reach the thermal equilibrium

# One-dimensional compressible flow

## Successive model simplifications

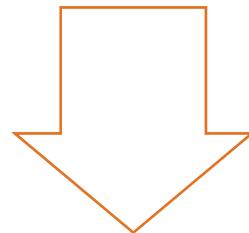
### Model 2

$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + (\rho v^2 + a_{\text{iso}}^2 \rho)_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

### Assumptions

- No diffusion
- One-dimensional flow
- No gravity
- Fluid is a polytropic, ideal gas
- Flow is isothermal



### Model 3

$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + a_{\text{iso}}^2 \rho_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

### Additional Assumption

- No convective acceleration
  - Valid for  $\text{Ma} < 0.3$

# One-dimensional compressible flow

## Successive model simplifications

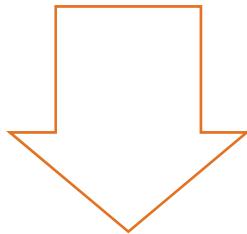
### Model 3

$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + a_{\text{iso}}^2 \rho_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

### Assumptions

- No diffusion
- One-dimensional flow
- No gravity
- Fluid is a polytropic, ideal gas
- Flow is isothermal
- No convective acceleration



### Model 4

$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + a_{\text{iso}}^2 \rho_z = -k_{\text{fric}} \frac{32\eta_0}{D^2} v, \quad k_{\text{fric}} \geq 1$$

### Additional Assumption

- Laminar flow
  - Valid for small Reynolds numbers
  - Effect of compressibility on the friction factor can be neglected

# One-dimensional compressible flow

## Successive model simplifications

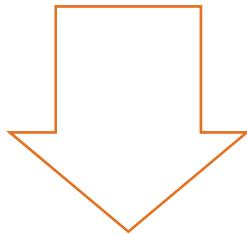
### Model 4

$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + a_{\text{iso}}^2 \rho_z = -k_{\text{fric}} \frac{32\eta_0}{D^2} v$$

### Assumptions

- No diffusion
- One-dimensional flow
- No gravity
- Fluid is a polytropic, ideal gas
- Flow is isothermal
- No convective acceleration
- Laminar flow



### Model 5

$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + a_{\text{iso}}^2 \rho_z = -k_{\text{fric}} \frac{32\eta_0}{D^2} \frac{1}{\rho_0} \rho v$$

### Additional Assumption

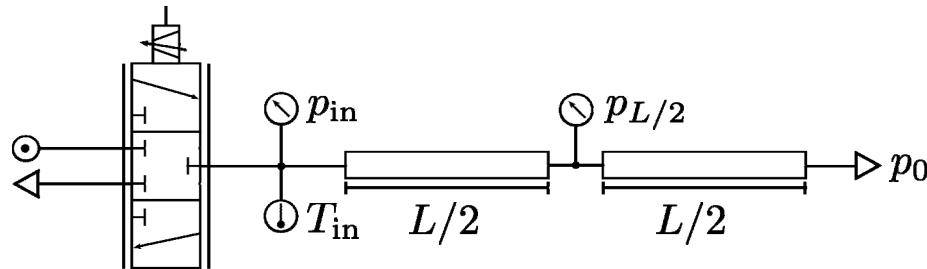
- Density is almost constant
  - Valid for small pressure and temperature changes

# Simulation and measurement results

## Testing scenarios

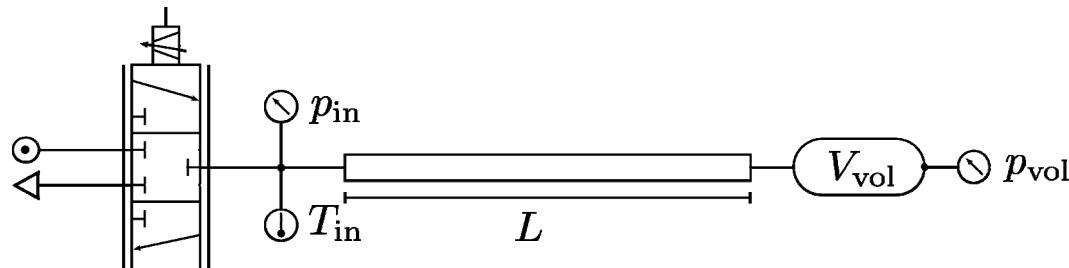
### Scenario 1

- Flow in ambient air



### Scenario 2

- Flow into a terminating volume



### Initial and boundary conditions

- Initial conditions  
 $p(z, 0) = p_0(z), \rho(z, 0) = \rho_0(z), v(z, 0) = 0$
- Left boundary condition  
 $p(0, t) = p_{\text{in}}(t), \rho(0, t) = \rho_{\text{in}}(t)$
- Right boundary condition scenario 1  
 $p(L, t) = p_0$
- Right boundary condition scenario 2  
$$p_{\text{vol}}(t) = \frac{m_{\text{vol}}(t)R_sT_{\text{vol}}(t)}{V_{\text{vol}}}$$
$$m_{\text{vol}}(t) = \int_0^t \dot{m}(L, \tau) d\tau + m_{\text{vol}}(0)$$
$$(c_{v,f,\text{vol}}m_{\text{vol}}T_{\text{vol}})_t(t) = (\dot{m}e)(L, t) + (pAv)(L, t) + \dot{Q}_{\text{vol}}(t)$$

# Simulation and measurement results

## Test bench

### Mechanical components

- Tube

Diameter	D	$8 \cdot 10^{-3} \text{ m}$
Length	L	19,83 m
Roughness	$\epsilon$	$1,5 \cdot 10^{-6} \text{ m}$

- Terminating volume

Volume	$V_{\text{vol}}$	$6,46 \cdot 10^{-4} \text{ m}^3$
Therm. restistance	$R_{\text{vol}}$	$4 \cdot 10^{-3} \text{ K/W}$

- Medium air

Ambient temperature	$T_0$	293,15 K
Ambient pressure	$p_0$	1,01 bar
Heat capacity ratio	$\gamma$	1,4

### Electrical components

- Pressure sensors
  - Festo pressure transmitter SPTE
- Ventil
  - Norgren proportional valve VP60
- Computer
  - Intel(R) Core(TM) i7-4770 CPU @ 3.40Ghz
  - Simulink Real-Time

### Implementation for simulation

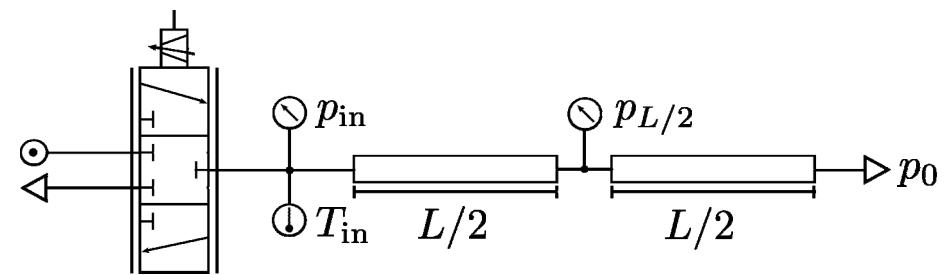
- The PDEs are discretized via the second-order finite difference method of MacCormack

# Simulation and measurement results

## Testing scenarios

### Scenario 1 – Flow into ambient air

- Step input signal
  - Opening the valve in 0,0175 s

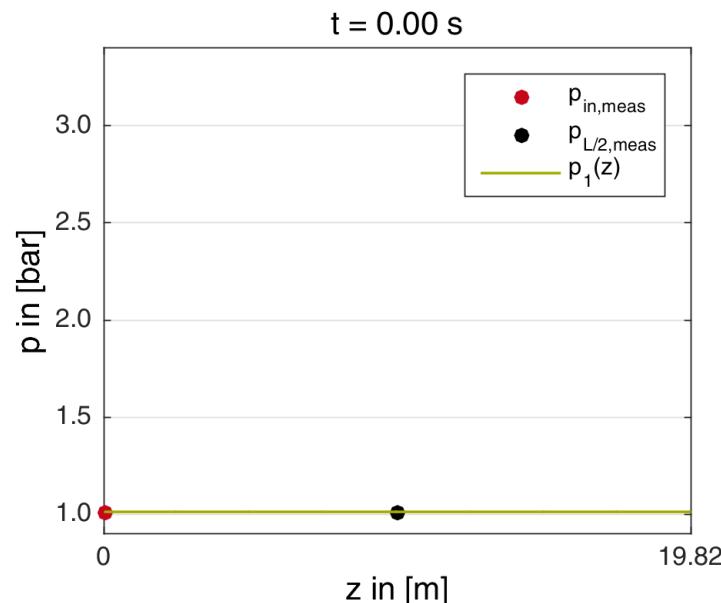


# Simulation and measurement results

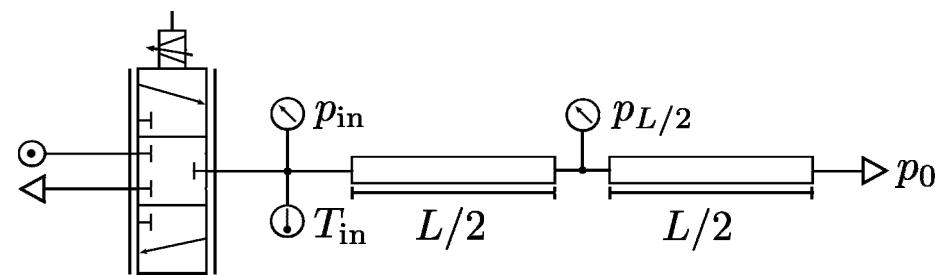
## Testing scenarios

### Scenario 1 – Flow into ambient air

- Step input signal
  - Opening the valve in 0,0175 s



### Shock wave



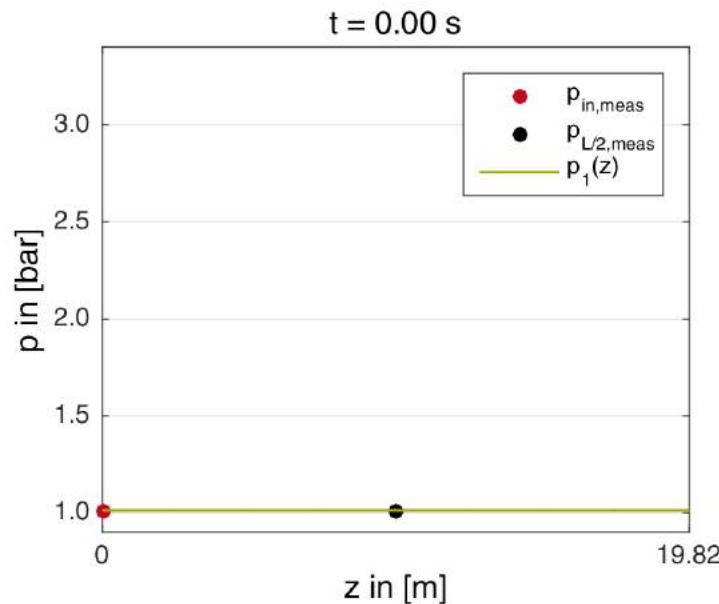
- Simulation of model 1
  - $v_{\max} = 111.06 \text{ m/s}$
  - $\text{Ma}_{\max} = 0.33$
  - $T_{\max} = 317.84 \text{ K}$

# Simulation and measurement results

## Testing scenarios

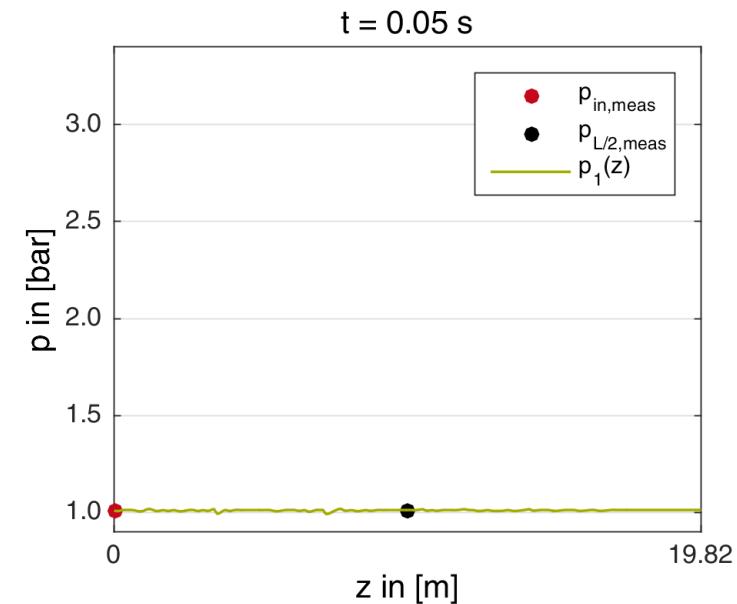
### Scenario 1 – Flow into ambient air

- Step input signal
  - Opening the valve in 0.0175 s



- Simulation of model 1
  - $v_{max} = 111,06 \text{ m/s}$
  - $Ma_{max} = 0,33$
  - $T_{max} = 317,84 \text{ K}$

### Shock wave

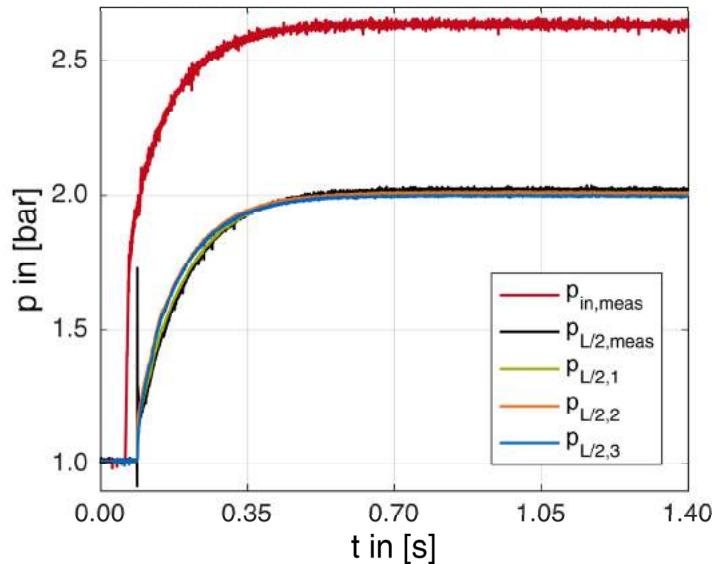


- The shock wave and its reflection are damped due to friction
- Scattering of the shock wave in the sensor causes measurement error
- CFL condition < 1 in the shock wave causes numerical dispersion

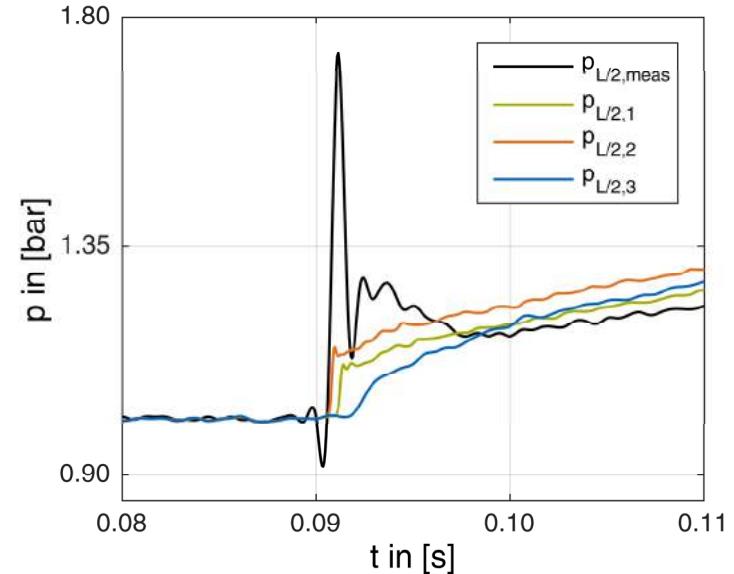
# Simulation and measurement results

## Testing scenarios

### Scenario 1 – Flow into ambient air



### Propagation speed



- Model 1

$$(\rho v)_t + (\rho v^2 + p)_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

$$(\rho e)_t + (v(\rho e + p))_z = \frac{1}{A} 2\pi \alpha r_i (T_0 - T) + f_{\text{comp}} \frac{\rho v^2 |v|}{2D}$$

- Model 2

$$(\rho v)_t + (\rho v^2 + a_{\text{iso}}^2 \rho)_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

- Model 3

$$(\rho v)_t + a_{\text{iso}}^2 \rho_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

- Model 1:  $v + \sqrt{\frac{\gamma p}{\rho}}$

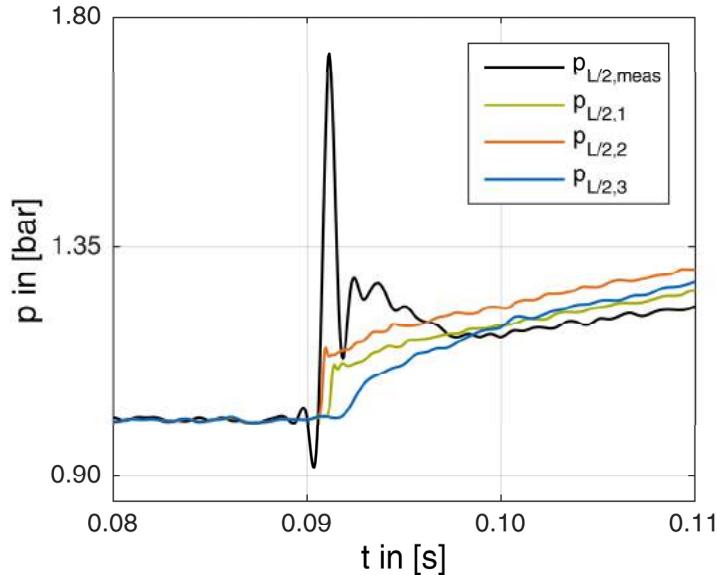
- Model 2:  $v + \sqrt{\gamma R_s T_0}$

- Model 3:  $\sqrt{\gamma R_s T_0}$

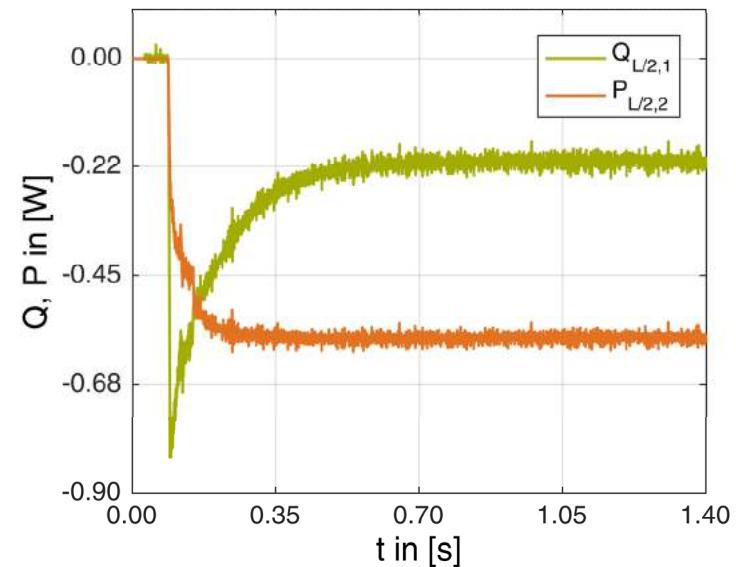
# Simulation and measurement results

## Testing scenarios

### Scenario 1 – Flow into ambient air



### Loss of energy



- Model 1

$$(\rho v)_t + (\rho v^2 + p)_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

$$(\rho e)_t + (v(\rho e + p))_z = \frac{1}{A} 2\pi \alpha r_i (T_0 - T) + f_{\text{comp}} \frac{\rho v^2 |v|}{2D}$$

- Model 2

$$(\rho v)_t + (\rho v^2 + a_{\text{iso}}^2 \rho)_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

- Model 3

$$(\rho v)_t + a_{\text{iso}}^2 \rho_z = -f_{\text{comp}} \frac{\rho v |v|}{2D}$$

At  $z = \frac{L}{2}$

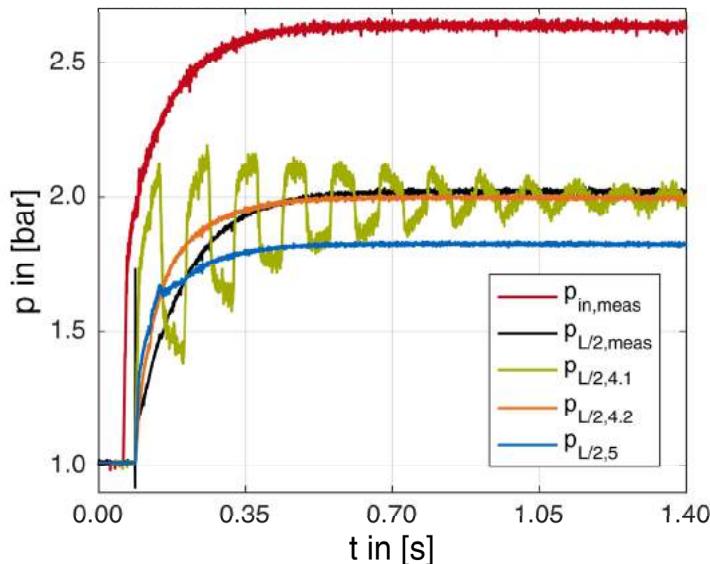
– Model 1:  $\dot{Q} = \frac{1}{A} 2\pi \alpha r_i (T_0 - T)$

– Model 2:  $P = f_{\text{comp}} \frac{\rho v^2 |v|}{2D}$

# Simulation and measurement results

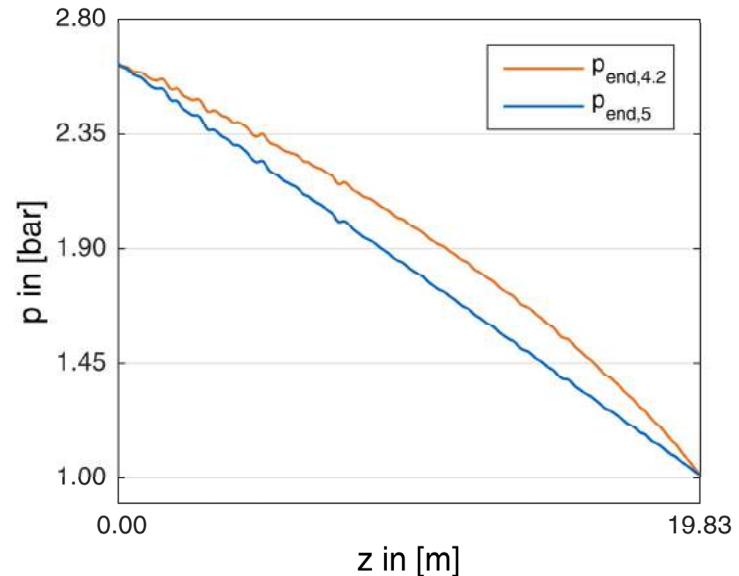
## Testing scenarios

### Scenario 1 – Flow into ambient air



- Model 4.1
$$(\rho v)_t + a_{iso}^2 \rho z = -k_{fric} \frac{32\eta_0}{D^2} v, \quad k_{fric} = 1$$
- Model 4.2
$$(\rho v)_t + a_{iso}^2 \rho z = -k_{fric} \frac{32\eta_0}{D^2} v, \quad k_{fric} = 15$$
- Model 5
$$(\rho v)_t + a_{iso}^2 \rho z = -k_{fric} \frac{32\eta_0}{D^2} \frac{1}{\rho_0} \rho v, \quad k_{fric} = 8$$

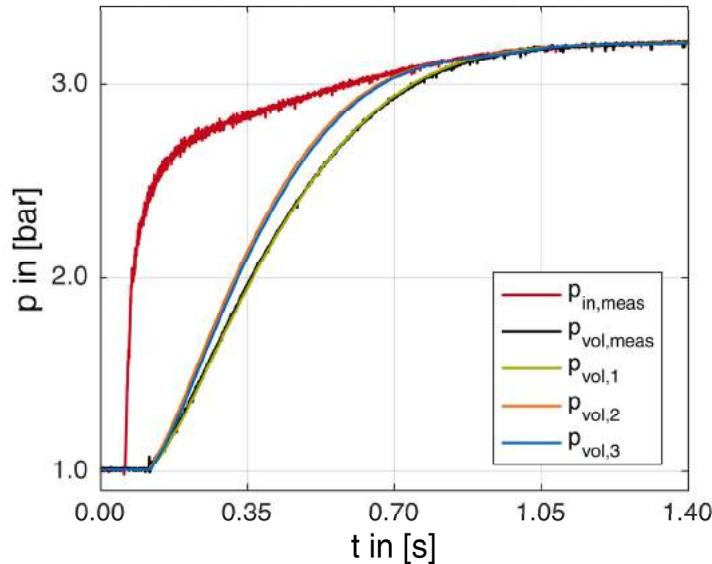
- Choice of  $k_{fric}$  for Model 5
  - Model 4: Friction laminar  
Proportional to  $v$
  - Model 5: Friction linear  
Proportional zu  $\rho v$
- Evident for the steady state solution
  - Pressure profile at  $t = 1.4$  s



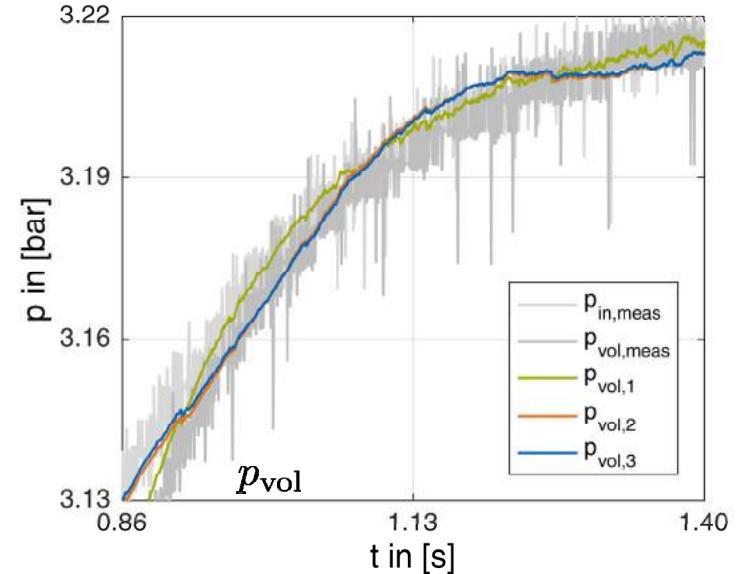
# Simulation and measurement results

## Testing scenarios

### Scenario 2 – Flow into a terminating volume



### Pressure in volume



- Model 1

$$(\rho v)_t + (\rho v^2 + p)_z = -f_{comp} \frac{\rho v |v|}{2D}$$

$$(\rho e)_t + (v(\rho e + p))_z = \frac{1}{A} 2\pi \alpha r_i (T_0 - T) + f_{comp} \frac{\rho v^2 |v|}{2D}$$

- Model 2

$$(\rho v)_t + (\rho v^2 + a_{iso}^2 \rho)_z = -f_{comp} \frac{\rho v |v|}{2D}$$

- Model 3

$$(\rho v)_t + a_{iso}^2 \rho_z = -f_{comp} \frac{\rho v |v|}{2D}$$

- Boundary conditions

- Model 1:

$$p_{vol} = \frac{\left( \int_0^t \dot{m}(\tau) d\tau + m_{vol} \right) R_s T_{vol}}{V_{vol}}$$

$$(c_{v,f,vol} m_{vol} T_{vol})_t = \dot{m}e + pAv + \dot{Q}_{vol}$$

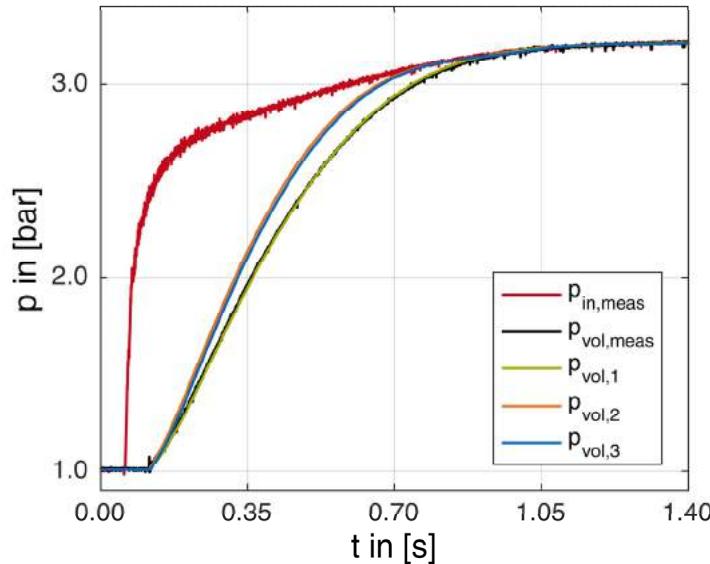
- Model 2 and 3

$$p_{vol} = \frac{\left( \int_0^t \dot{m}(\tau) d\tau + m_{vol} \right) R_s T_0}{V_{vol}}$$

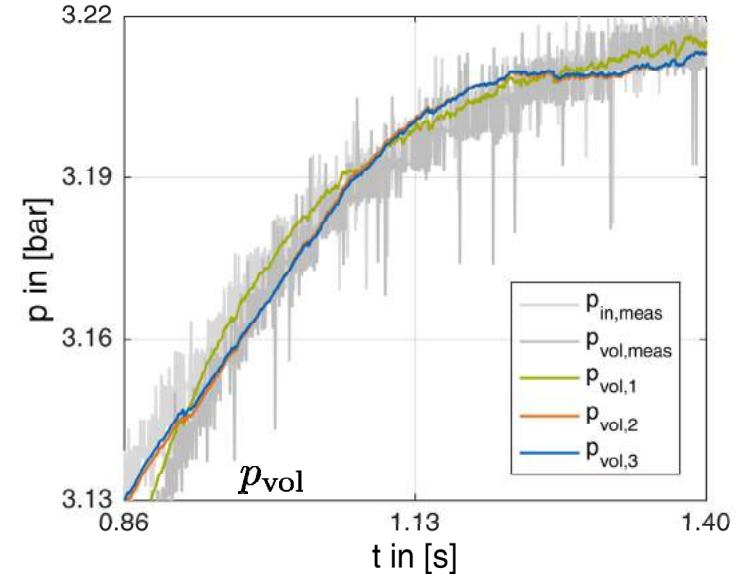
# Simulation and measurement results

## Testing scenarios

### Scenario 2 – Flow into a terminating volume



### Pressure in volume



- Deviation from measurement at  $t = 0.5$  s
  - Model 1: 0,45 %
  - Model 2: 8,02 %
- Total energy loss
  - Ratio model 1 to model 2: 1.36

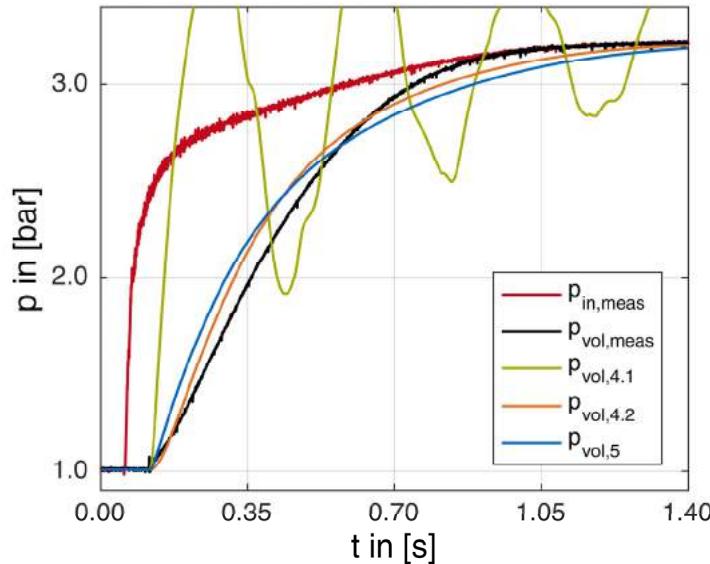
### Boundary conditions

- Model 1:
$$p_{vol} = \frac{\left( \int_0^t \dot{m}(\tau) d\tau + m_{vol} \right) R_s T_{vol}}{V_{vol}}$$
$$(c_{v,f,vol} m_{vol} T_{vol})_t = \dot{m}e + pAv + \dot{Q}_{vol}$$
- Model 2 and 3:
$$p_{vol} = \frac{\left( \int_0^t \dot{m}(\tau) d\tau + m_{vol} \right) R_s T_0}{V_{vol}}$$

# Simulation and measurement results

## Testing scenarios

### Scenario 2 – Flow into a terminating volume



- Model 4.1
$$(\rho v)_t + a_{iso}^2 \rho z = -k_{fric} \frac{32\eta_0}{D^2} v, \quad k_{fric} = 1$$
- Model 4.2
$$(\rho v)_t + a_{iso}^2 \rho z = -k_{fric} \frac{32\eta_0}{D^2} v, \quad k_{fric} = 15$$
- Model 5
$$(\rho v)_t + a_{iso}^2 \rho z = -k_{fric} \frac{32\eta_0}{D^2} \frac{1}{\rho_0} \rho v, \quad k_{fric} = 8$$

# Summary and outlook

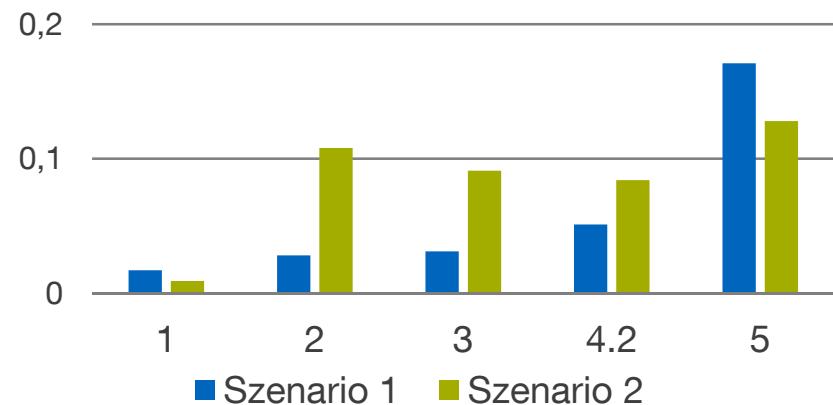
## Summary

### Assumptions and effects

- Basic assumptions – Model 1
- Isothermal flow – Model 2
  - Two instead of three equations
  - Conservation of energy is violated
- No convective acceleration – Model 3
  - Coefficients of the derivatives are constant (equations are semilinear instead of quasilinear)
  - Wave propagation speed is incorrect
- Laminar flow – Model 4
  - Correlations are not necessary
  - Effect of friction without amplification factor is underestimated
- Constant density – Model 5
  - Linear equations
  - Relatively severe error due to the approximations

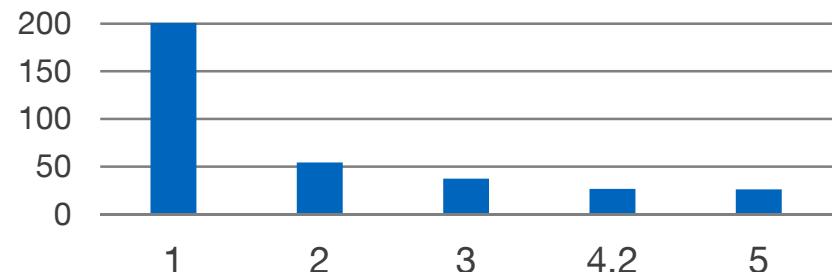
### Root mean square

- Difference between simulated and measured pressure (in bar)



### Computational time

- Simulation of scenario 1 (in s) with  $N_z = 992$  and  $N_t = 24\,270 - 30\,514$



# Summary and outlook

## Outlook

### Applications of the models

- Complex models for
  - Simulation
  - Optimization
  - Distributed-parameter feedforward control
  - Lumped-parameter feedback control
- Less complex models
  - If pressure changes are relatively small
  - Suited if the system exhibits a terminating volume at the end of the transmission line (e.g. piston)
  - Time-critical applications
  - Distributed-parameter feedback control



# References

## Literature

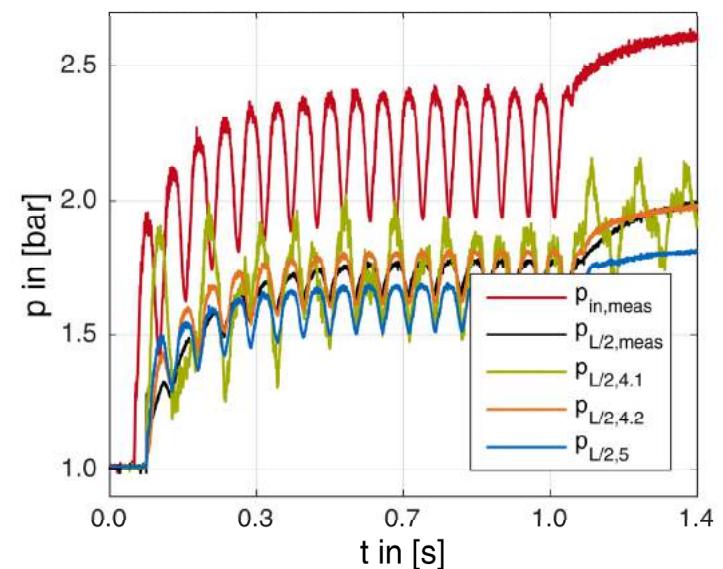
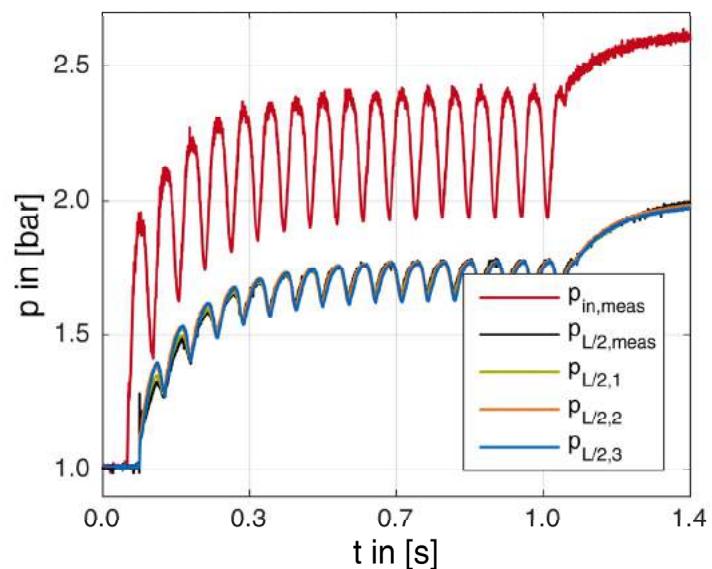
- LeVeque, R. J., Numerical Methods for Conservation Laws. Birkhäuser. 1999.
- Toro, E. F., Riemann Solvers and Numerical Methods for Fluid Dynamics: A Practical Introduction. Springer. 2009.
- Munson, B. R., Fundamentals of Fluid Mechanics. Wiley. 2013.
- Krichel, S. V.; Sawodny, O.: Non-linear friction modelling and simulation of long pneumatic transmission lines. In: Math. Comput. Model. Dyn. Syst. 20 (2013), Nr. 1, S. 23–44.
- Rager, D.; Neumann, R.; Murrenhoff, H.: Simplified fluid transmission line model for pneumatic control applications. In: Proc. 14th Scandinavian International Conference on Fluid Power (SICFP15). Tampere, Finland, 2015.
- Stecki, J. S.; Davis, D. C.: Fluid transmission lines—distributed parameter models part 1: A review of the state of the art. In: Proc. IME J. Power Energ. 200 (1986), Nr. 41, S. 215–228.

# Sinusanregung

## Testing scenarios

### Scenario 1

- Excitation with maximum frequency ( $\sim 16$  Hz)

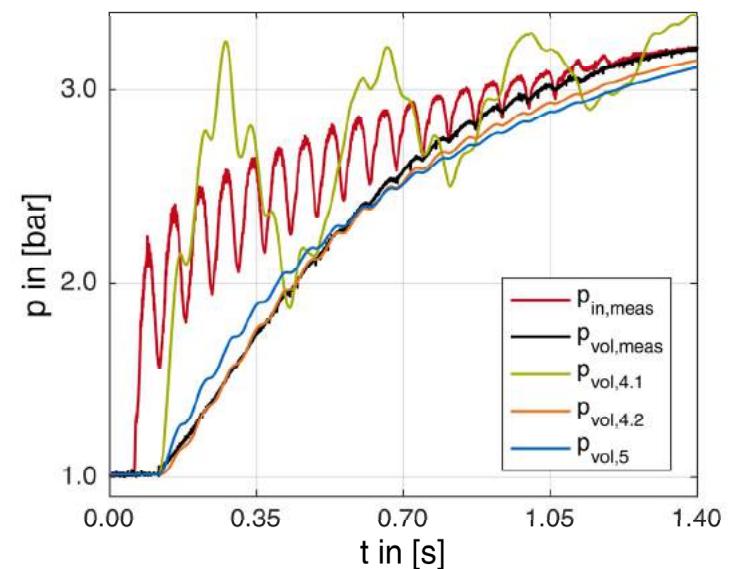
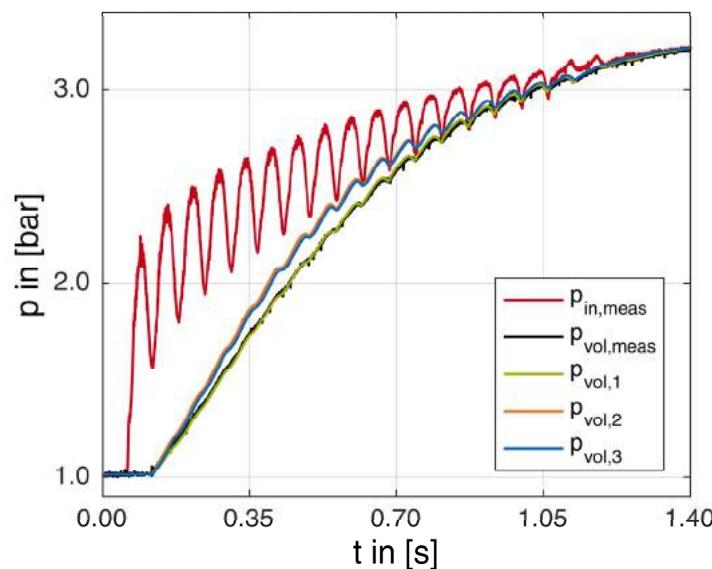


# Sinusanregung

## Testing scenarios

### Scenario 2

- Excitation with maximum frequency (~ 16 Hz)



# Derivation of Model 3

## Linearization of the pressure function

- Formulierung der Eulergleichungen durch Entropieerhaltung

$$\rho_t + (\rho v)_z = 0$$

$$(\rho v)_t + (\rho v^2 + p)_z = 0$$

$$S_t + vS_z = 0$$

- Definition der Entropie

$$S = c_v \log \left( \frac{p}{\rho^\gamma} \right) + \kappa$$

- Auflösen nach Druck

$$p(\rho) = \exp \left( \frac{S}{c_v} + \kappa \right) \rho^\gamma$$

- Taylorentwicklung für Funktion des Drucks

$$\begin{aligned} p(\rho) &= p(\rho_0) + \gamma \exp \left( \frac{S}{c_v} + \kappa \right) \rho_0^{\gamma-1} (\rho - \rho_0) + \dots \\ &= p(\rho_0) + \gamma \frac{p(\rho_0)}{\rho_0} (\rho - \rho_0) + \dots \end{aligned}$$

- Assumption von kleinen Änderungen

$$p(\rho) = p(\rho_0) + \gamma R_s T_0 (\rho - \rho_0)$$

- Isotherme Eulergleichungen

$$\rho_t + (\rho v)_z = 0$$

$$(\rho v^2 + a_{\text{iso}}^2 \rho)_z = 0 \quad \text{with} \quad a_{\text{iso}} = \sqrt{\gamma R_s T_0}$$